

## ADDITIONAL MONOTONICITY PROPERTIES AND INEQUALITIES FOR THE ZEROS OF BESSEL FUNCTIONS (\*)

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### 1. INTRODUCTION.

For  $\nu \geq 0$  we denote with  $j_{\nu k}$  and  $c_{\nu k}$  the  $k$ -th positive zeros of the Bessel functions  $J_\nu(x)$  of the first kind and of the general cylinder function

$$C_\nu(x; \gamma) = C_\nu(x) =: \cos \gamma J_\nu(x) - \sin \gamma Y_\nu(x) \quad 0 \leq \gamma < \pi$$

where  $Y_\nu(x)$  is the Bessel function of the second kind.

In [2] A. Elbert and A. Laforgia introduced the notation  $j_{\nu \kappa} = c_{\nu k}$  where  $\kappa = k - \gamma/\pi$  and  $k-1 < \kappa < k$ . When  $\kappa = k$  we get the zeros of the function  $J_\nu(x)$ . This notation has been used to prove several *monotonicity, concavity, convexity* properties of  $j_{\nu \kappa}$  as a function of  $\nu$  for  $\kappa$  fixed [2,3,4,5,9]. In particular we know that for  $\nu \geq 0$   $j_{\nu \kappa}$  is concave for  $\kappa \geq 0.344\dots$  and  $j_{\nu \kappa}^2$  is convex for  $\kappa \geq 0.7070\dots$  [5].

We observe that the study of the properties of  $j_{\nu \kappa}$  was originated by the paper [19] of Putterman, Kac and Uhlenbeck. They have proposed a quantum mechanical explanation for the origin of the vortex lines produced in superfluid Helium when its container is rotated.

J.T. Lewis and M.E. Muldoon [16] have proved some monotonicity results using the Hellmann-Feynmann theorem of quantum chemistry [8,10,11] that here we recall

**Hellmann-Feynmann Theorem.** Let's be a pseudo inner product space with a pseudo-inner product  $\langle \cdot, \cdot \rangle$ . Let  $(H_\nu)$  be a family of symmetric operators on an inner product space and for  $\nu \in (a,b)$  let  $\psi_\nu$  be an eigenvector (eigenfunction) of  $H_\nu$  corresponding to an eigenvalue  $\lambda_\nu$ .

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(\*) Work sponsored by Ministero dell'Università e della Ricerca Scientifica e Tecnologica of Italy.

Suppose that for  $\mu \in (a, b)$

$$\langle \Psi_\mu, \Psi_\nu \rangle \rightarrow \langle \Psi_\mu, \Psi_\mu \rangle \neq 0$$

as  $\mu \rightarrow \nu$  and that

$$\lim_{\mu \rightarrow \nu} \left\langle \frac{H_\mu - H_\nu}{\mu - \nu}, \Psi_\mu, \Psi_\nu \right\rangle$$

exists.

Moreover we define

$$\left\langle \frac{\partial H_\nu}{\partial \nu}, \Psi_\mu, \Psi_\nu \right\rangle := \lim_{\mu \rightarrow \nu} \left\langle \frac{H_\mu - H_\nu}{\mu - \nu}, \Psi_\mu, \Psi_\nu \right\rangle$$

Then

$$\frac{d\lambda_\nu}{d\nu} = \frac{\left\langle \frac{\partial H_\nu}{\partial \nu}, \Psi_\mu, \Psi_\nu \right\rangle}{\langle \Psi_\mu, \Psi_\nu \rangle}$$

This version of the Theorem is taken from [11]. As a consequence of the Hellmann- Feynmann Theorem, Lewis and Muldoon [16] have proved that  $j_{\nu 1}/\nu$  decreases with  $\nu > 0$  and that  $j_{\nu 1}^2/\nu$  increases with  $\nu$ , ( $3 \leq \nu < \infty$ ).

More recently C. Giordano and L.G. Rodonò proved that for  $\nu > 0$  and  $\kappa \geq 0.7070\dots$  there exists a value  $\nu_\kappa$  such that the function  $j_{\nu \kappa}^2/\nu$  decreases on  $(0, \nu_\kappa)$  and increases on  $(\nu_\kappa, \infty)$ . Furthermore they showed that  $j_{\nu \kappa}^2/\nu$  is convex on  $(0, \nu_\kappa)$  [9]. We also recall that M.E.H. Ismail and M.E. Muldoon [12] proved the stronger monotonicity result that  $j_{\nu 1}^2/(\nu+1)$  is an increasing function of  $\nu$  on  $(-1, \infty)$ .

Further properties for  $j_{\nu \kappa}$  and  $c_{\nu \kappa}$  can be established by the integral Watson formula [22, p. 508]

$$(1.1) \quad \frac{d c_{\nu \kappa}}{d\nu} = 2c_{\nu \kappa} \int_0^\infty K_0(2c_{\nu \kappa} \sinh t) e^{-2\nu t} dt, \quad \kappa=1,2 \dots$$

where  $K_0(x)$  is the standard modified Bessel function.

In section 2 we recall some monotonicity properties recently established by A.Laforgia, M.E.Muldoon and L.Lorch and in section 3 we examine their application to the zeros of generalized Airy functions [6,15,17].

## 2. MONOTONICITY RESULTS.

In [17] L.Lorch has deduced the inequality

$$\frac{\nu}{c_{\nu \kappa}} \cdot \frac{d c_{\nu \kappa}}{d\nu} < 1, \quad \text{for } c_{\nu \kappa} \geq \nu > 0,$$

which implies that, for  $k=1,2,\dots$ ,  $0 < \nu < \infty$

$$\frac{\nu}{c_{\nu k}} \quad \text{increases for } c_{\nu k} > \nu .$$

When  $c_{\nu k} > \nu + \pi/4$ , there follows the stronger result that [14,17]

$$\frac{\nu + \frac{1}{2}}{c_{\nu k}} \quad \text{increases for } 0 \leq \nu < \infty .$$

More generally A.Laforgia and M.E.Muldoon in [14] have showed the inequality

$$c_{\nu k} > \nu + c_{0k} , \quad 0 < \nu < \infty , \quad k=2,3,\dots$$

and

$$(2.1) \quad (\nu + a_k) \frac{d c_{\nu k}}{d \nu} < c_{\nu k} , \quad 0 < \nu < \infty , \quad k=2,3,\dots$$

where the positive number  $a_k$  is defined by

$$a_k = c_{0k}^{-1} - c_{0k}^{-1} \left[ \frac{d c_{\nu k}}{d \nu} \right]_{\nu=0}$$

The previous results hold also in the case  $k=1$ , but  $0 \leq \nu \leq \pi/2$ .

Inequality (2.1) can be employed to yield the stronger monotonicity result

$$\frac{\nu + a_k}{c_{\nu k}} \quad \text{increases for } \nu > 0 \quad \text{when } k=2,3,\dots$$

If  $0 \leq \nu \leq \pi/2$  this monotonicity holds also for  $k=1$ .

Differentiating we get

$$c_{\nu k} \left( \frac{\nu + a_k}{c_{\nu k}} \right)' = 1 - \frac{\nu + a_k}{c_{\nu k}} \frac{d c_{\nu k}}{d \nu} > 0$$

The previous results can be employed to deduce other monotonicity results as the following ones recently found by A.Laforgia, M.E.Muldoon and L.Lorch [15,17].

We recall now some of them.

If  $\nu > 0$  the function

$$\left( \frac{c_{\nu k}}{\nu} \right)^{\nu} \quad \text{increases for } c_{\nu k} > \nu + \frac{\pi}{4}$$

and

$$\left( \frac{c_{\nu k}}{2\nu} \right)^{2\nu} \quad \text{decreases for } 0 < \nu \leq c_{\nu k} \leq 2\nu$$

The proof is based on the properties [17] of the function

$$\delta_{\nu} = \frac{d}{d \nu} \left\{ \ln \left[ \frac{c_{\nu k}}{\beta \nu} \right]^{\nu} \right\} , \quad \nu > 0 , \quad \beta > 0$$

in the case  $\beta=1, \beta=2$ , respectively.

The monotonicity of  $\delta_\nu$  implies that if  $\delta_\nu > 0$ , then  $[c_{\nu k}/(2\nu)]^{2\nu}$  increases for  $0 < \nu \leq \mu$ , while  $\delta_\nu \leq 0$  implies that this function of  $\nu$  decreases for all  $\nu > \mu$ , provided of course that  $c_{\nu k} > \nu + \pi/4$ .

The hypothesis that  $c_{\nu k} > \nu + \pi/4$  is satisfied when  $k=2,3,\dots$  and for any  $c_{\nu 1}$  for which  $0 \leq \gamma \leq \pi/2$ ; in particular for  $c_{\nu 1} = j_{\nu 1}$ .

In the case  $k=1$ ,  $c_{\nu k} = j_{\nu 1}$  there exists a value  $\nu = \nu_1$  such that

$$(2.2) \quad \left(\frac{j_{\nu 1}}{2\nu}\right)^{2\nu} \quad \text{increases,} \quad 0 < \nu \leq \nu_1$$

and

$$(2.3) \quad \left(\frac{j_{\nu 1}}{2\nu}\right)^{2\nu} \quad \text{decreases,} \quad \nu \geq \nu_1$$

M.E.Muldoon has calculated some values of  $[j_{\nu 1}/(2\nu)]^{2\nu}$  near 1 from which it derives that  $1.003 < \nu_1 < 1.006$ .

If the order  $\nu$  is kept constant, but  $k$  or  $\gamma$  varies in  $c_{\nu k}(\gamma)$  it could be convenient to use the notation introduced by Á.Elbert and A.Laforgia in [2]. They show that  $j_{\nu k}$  increases with  $k > 0$ , for fixed  $\nu$ .

L.Lorch in [17] extends (2.2), (2.3) by implying the existence of unique  $\nu_\kappa$  such that

$$(2.4) \quad \left(\frac{j_{\nu \kappa}}{2\nu}\right)^{2\nu} \quad \text{increases for } 0 < \nu \leq \nu_\kappa \text{ and decreases for } \nu \geq \nu_\kappa,$$

where  $\nu_\kappa$  is an increasing function of  $\kappa > 0$ .

The behaviour of  $[c_{\nu k}/(2\nu)]^{2\nu}$  suggests to consider the character of  $[c_{\nu k}/(2\nu)]^{2\nu}/\nu$  [15,17].

First of all A.Laforgia and M.E.Muldoon in [15] have obtained that for some  $0 < \nu_0 \leq 1/2$  and some  $k=1,2,\dots$ , under some hypothesis, the function  $[c_{\nu k}/(2\nu)]^{2\nu}/\nu$  decreases as  $\nu$  increases for  $0 < \nu \leq \nu_0$ .

Recently L.Lorch has established in [17] that  $[c_{\nu 1}/(2\nu)]^{2\nu}/\nu$  decreases for  $0 < \nu < \infty$  if  $c_{\nu 1} > \nu$ . The proof is divided into four parts. In each subinterval  $\lambda < \nu \leq \mu$  he has proved that the expression

$$\frac{1}{2} \frac{d}{d\nu} \left\{ \ln \left[ \left(\frac{c_{\nu 1}}{2\nu}\right)^{2\nu} \cdot \frac{1}{\nu} \right] \right\}$$

is negative, using previous results in [15, 17].

### 3. APPLICATION TO THE ZEROS OF GENERALIZED AIRY FUNCTIONS.

The generalized Airy functions are the solutions on  $0 \leq x < \infty$  of the equation

$$(3.1) \quad y'' + x^\alpha y = 0$$

where  $\alpha$  is a positive number. The Airy functions correspond to the case  $\alpha = 1$ .

The case  $\alpha \neq 1$  arises in a fundamental way in the asymptotic solution of certain kinds of differential equations [18,20,21]. A function of this kind also occurs in work concerned with weighted averages of a function at a jump discontinuity [1].

Equation (3.1) is connected to the Bessel equation

$$t^2 \frac{d^2 u}{dt^2} + t \frac{du}{dt} + (t^2 - \nu^2) u = 0$$

by means of the transformations

$$(3.2) \quad y(x) = x^{\nu/2} u(t) \quad t = 2\nu x^{1/(2\nu)}$$

where  $\nu = 1/(\alpha + 2)$ .

So it is possible to use many known results about monotonicity with respect to order of zeros of Bessel functions. We denote  $a_{\alpha k}$  the  $k$ -th positive zero of a solution of (3.1).

Because of (3.2), it is clear that

$$(3.3) \quad a_{\alpha k} = \left( \frac{c_{\nu k}}{2\nu} \right)^{2\nu} \quad \text{where } 0 < \nu = \frac{1}{\alpha + 2}.$$

Using the connection (3.3) between zeros of cylinder functions and generalized Airy functions, it is possible to obtain results for  $a_{\alpha k}$ . As Corollary to Theorem 2.1 in [14] there follows that for each  $k=2,3,\dots$   $a_{\alpha k}$  decreases to 1 as  $\alpha$  increases,  $0 < \alpha < \infty$ . If  $y(0) = 0$  this decrease holds also for  $a_{\alpha 1}$ .

Muldoon's bounds for  $\nu_1$  permit to assert (2.2) and (2.3) with  $1.003 < \nu_1 < 1.006$ .

Consequently

$$(3.4) \quad \begin{array}{ll} a_{\alpha 1} & \text{increases } -1.00596 < -2 + 1/\nu_1 \leq \alpha < \infty \\ a_{\alpha 1} & \text{decreases } -2 < \alpha \leq -2 + 1/\nu_1 < -1.00299 \end{array}$$

Taking into account (2.4), the result (3.4) can be extended to

$$(3.5) \quad \begin{array}{ll} a_{\alpha k} & \text{increases } -2 + 1/\nu_k \leq \alpha < \infty \\ a_{\alpha k} & \text{decreases } \alpha < -2 + 1/\nu_k \end{array}$$

This implies that

$(\alpha + 2) a_{\alpha 1}$  increases,  $-1.0029 < \alpha < \infty$

and more generally

$(\alpha + 2) a_{\alpha k}$  increases,  $-2 + 1/\nu_k < \alpha < \infty$

for  $k = 1, 2, \dots$ .

In [6] Á. Elbert and A. Laforgia continued the investigation about the behaviour of  $a_{\alpha k}$  and they proved the following result

**Theorem.** For each fixed  $k = 2, 3, \dots$ , let  $a_{\alpha k}$  be the  $k$ -th positive zero of (3.1). Then for  $\alpha \geq 0$ , the function  $\log a_{\alpha k}$  is convex.

The first important consequence of this theorem is the convexity of  $a_{\alpha k}$ , since a positive log-convex function is also convex.

The proof is based on some known properties of  $c_{\nu k}$  and the Watson's formula (1.1).

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